

Normal mode analysis, fully discrete case

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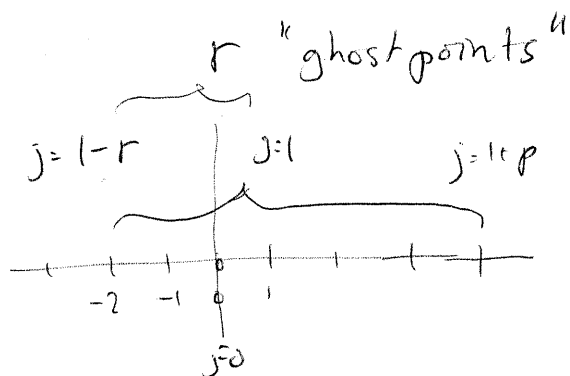
Analysis similar to the semi-discrete case with the changes "Laplace transform" \rightarrow "Z-transform" and " $\text{Re } s \geq 0$ " \rightarrow " $|z| \geq 1$ ".

Consider constant coefficient, 1/4-plane problems:

$$(*) \begin{cases} u_j^{n+1} = Q u_j^n + k F_j^n & j=1, 2, \dots \\ B u^n = g^n, \quad \|u\| < \infty \\ u_j^0 = f_j \end{cases}$$

where

$$Q = \sum_{m=-r}^p a_m E^m$$



$a_m = a_m(h, k)$ bounded for k/h fixed.

$$B u^n \sim u_v^n = \sum_{j=1}^q b_{vj} u_j^n + g_v^n + \sum_{j=1}^q b'_{vj} u_j^{n-1} + \dots$$

$v = 0, \dots, r+1$

May also contain different time levels $(n, n-1, \dots)$

Ex:

$$\begin{cases} u_j^{n+1} = \frac{u_{j-1}^n + u_{j+1}^n}{2} + \frac{k}{h} \frac{(u_{j+1}^n - u_{j-1}^n)}{2} & \text{Lax-Friedrichs} \\ u_0^n - 2u_1^n + u_2^n = 0 & (\text{or } u_0^n = g^n) \\ u_j^0 = f_j \end{cases}$$

for $u_t = u_x$.

As in semidiscrete case, we analyze:

(2)

a) $g=0, f \neq 0, F=0$

b) $g \neq 0, f=0, F=0$.

\Rightarrow Full solution is a sum of these.

Case a) $g=0, f \neq 0$.

Suppose $Q\phi_j = z\phi_j, B\phi_j = 0, \|\phi_j\| < \infty$

\Rightarrow Then

$$u_j^n = z^n \phi_j$$

is a solution of (*) (check: $Qu_j^n = z^n Q\phi_j = z^{n+1} \phi_j = u_j^{n+1}$
+ BC is ok.)
with initial data $f_j = \phi_j$.

$$\Rightarrow \|u^n\| = |z|^n \|\phi_j\| = |z|^{t_n/k} \|\phi_j\|$$

This cannot be bounded by $Ke^{\alpha t_n}$ independent of k if $|z| > 1$.

Godunov Rya benkii condition for fully discrete case:

If z is an eigenvalue of (Q, B)

then $|z| \leq 1$ is necessary for stability.

Case b) $g \neq 0, f = 0.$

(3)

(unilateral)

Use Z-transform for a sequence $\{u_n\}_g$

$$\hat{u}(z) = \sum_{n=0}^{\infty} z^{-n} u_n \quad (|z| \text{ large enough to get convergence})$$

Since $u_n \sim e^{\alpha n} = (e^{\alpha})^n$ we need $|z| > e^{\alpha}$

Then:

$$1) u^n = \frac{1}{2\pi i} \oint_C \hat{u}(z) z^{n-1} dz$$

(inversion)

$C = \{|z| = R\}$ with R large enough.

$$2) \widehat{u^{n+1}} = z \hat{u}(z) - z \cdot u^0$$

($> e^{\alpha k}$)

$$3) \sum |u_n|^2 = \frac{1}{2\pi} \int_{|z|=1} |\hat{u}(z)|^2 dz \quad (\text{Parseval})$$

Z-transform of (*) with $f = F = 0$:

$$(a) z \hat{u}_j(z) = Q \hat{u}_j(z)$$

$$\left\{ \begin{array}{l} \hat{B} \hat{u}_j(z) = \hat{g}(z) \\ (c) \end{array} \right., \quad \| \hat{u}_j(z) \| < \infty \quad (b)$$

Same eigenvalue problem as before!

Need to be able to solve it with $|z|$ large enough, to use inversion formula.

We can solve this in the same way as in semi-discrete case:

1) (a) is a difference equation parameterized by z .

If $Q = \sum_{m=-r}^p a_m E^m$, then the general solution

$$\hat{u}_j(z) = \sum_{m=1}^{r+p} \sigma_m k_m^j(z)$$

where $k_m(z)$ are roots of the characteristic polynomial:

$$a_{-r} + k a_{-r+1} + \dots + k^r (a_0 - z) + \dots + k^{r+p} a_p = P_z(k)$$

σ_m are free parameters.

2) Boundary cond. (b) $\|\hat{u}\| < \infty$ (4)

Need $\sigma_m = 0$ for $|k_m| \geq 1$ to satisfy (b).

⊙ Theorem: If von Neumann cond. satisfied for periodic problem, then:

- 1) No roots k_m with $|k_m| = 1$ if $|z| > 1$.
- 2) Precisely r roots with $|k_m| < 1$ if $|z| > 1$.

⇒ For $|z| > 1$ the solution is of the form

(***)
$$\hat{u}_j(z) = \sum_{m=1}^r \sigma_m k_m^j(z)$$

← note r

if periodic problem is stable.

3) Boundary condition (c) $\vec{B}\vec{u} = \vec{g}$.

Enter the form (***) into boundary conditions:

$$a_n u_\nu = \sum b_{\nu j} u_j^n + \sum \hat{b}_{\nu j} u_j^{n-1} + g_\nu \quad \nu = 0, \dots, r+1$$

$$\Rightarrow \hat{u}_\nu(z) = \sum (b_{\nu j} \hat{u}_j(z) + \hat{b}_{\nu j} \frac{1}{z} \hat{u}_j(z)) + g_\nu$$

$$\Rightarrow C(z) \begin{pmatrix} \sigma_1 \\ \vdots \\ \sigma_r \end{pmatrix} = \begin{pmatrix} g_0 \\ \vdots \\ g_{r+1} \end{pmatrix} \quad \text{where } C \in \mathbb{C}^{r \times r}$$

Same kind of (2) as before.

Need $\det C(z) \neq 0$ for $|z| > 1$. Otherwise there is a solution to eigenvalue problem with $\sigma_j = 0$.

∴ GR-condition $\Leftrightarrow \det C(z) \neq 0$ for $|z| > 1$.

Kreiss condition:

$$\det C(z) \neq 0 \text{ for } |z| \geq 1.$$

(Equivalent:

$$|\hat{u}_j(z)| \leq K |\hat{g}(z)| \quad (|z| \geq 1, K \text{ indep of } z.)$$

Together with energy estimate for full Cauchy problem ($j = -\infty \dots \infty$) and technical conditions it is sufficient for strong stability. $\|u^{n+1}\| \leq \|u^n\|$ = stronger than stability !! (cf Theorem 2.9 in book.)

Note $C(z)$ is defined with the roots k_m such that $|k_m| < 1$ for $|z| > 1$. When $|z| = 1$ k_m may have $|k_m| = 1$ and there may be more roots $|k_{m'}(z)| = 1$. These latter ones are not part of $C(z)$.

Example

⑥

Lax - Friedrichs scheme:

$$\begin{cases} u_j^{n+1} = \frac{u_{j+1}^n + u_{j-1}^n}{2} + \frac{\lambda}{2} (u_{j+1}^n - u_{j-1}^n) & j=1, \dots \\ u_2 - 2u_1 + u_0 = g \end{cases}$$

⊗ Z-transform \Rightarrow

$$(1) \quad z \hat{u}_j = \frac{\lambda+1}{2} \hat{u}_{j+1} + \frac{-\lambda+1}{2} \hat{u}_{j-1}$$

Characteristic polynomial =

$$(\lambda+1)k^2 - 2zk + (1-\lambda) = 0$$

$$\Rightarrow k_{1,2} = \frac{z \pm \sqrt{z^2 + \lambda^2 - 1}}{1 + \lambda}$$

Pick $|k| < 1$ for $|z| > 1$ (always possible by theorem since L-F stable)

$$k_1 = \frac{z - \sqrt{z^2 + \lambda^2 - 1}}{1 + \lambda}$$

$\therefore \hat{u}_j = \sigma_1 k_1^j(z)$ satisfies (1) and $\|\hat{u}_j\| < \infty$ for $|z| > 1$.

Boundary conditions:

$$\hat{u}_2 - 2\hat{u}_1 + \hat{u}_0 = \hat{g} \Rightarrow$$

$$\sigma(k_1^2 - 2k_1 + 1) = \hat{g} \Rightarrow \sigma(k_1 - 1)^2 = \hat{g}$$

$$C(z) = (K(z) - 1)^2 \quad (\text{Scalar here } \textcircled{7} \text{ since } r=1)$$

$$\det C(z) = 0 \Leftrightarrow K(z) = 1.$$

Not possible for $|z| > 1$ by construction ($|K| < 1$)

How about for $|z| = 1$?

$$K=1 \Rightarrow (\lambda+1)z^2 - 2z \cdot 1 + (1-\lambda) = 0$$

$\Rightarrow z=1$ only possibility

$$\text{But } K_1(z) = \frac{1 - \sqrt{1 + \lambda^2 - 1}}{1 + \lambda} = \frac{1 - \lambda}{1 + \lambda} < 1.$$

(i.e. $K_2(1) = 1$
not K_1 .)

$\therefore \det C(z) \neq 0$ for $|z| \geq 1$. Kreiss condition is satisfied.

Remarks

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- Multistep methods can be transformed to one-step methods for systems as before.

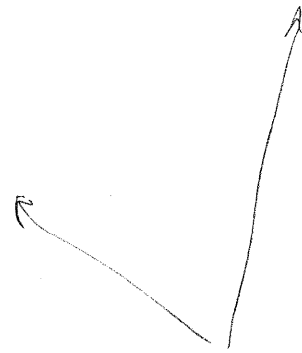
$$u^{n+1} = Q_0 u^n + Q_1 u^{n-1} + \dots + Q_s u^{n-s}$$

$$\text{Let } v = \begin{pmatrix} u^n \\ \vdots \\ u^{n-s} \end{pmatrix}$$

$$\text{Then } v^{n+1} = \begin{pmatrix} Q_0 & Q_1 & \dots & Q_s \\ \mathbf{I} & 0 & & 0 \\ & \mathbf{I} & & \\ & & \ddots & \\ & & & \mathbf{I} & 0 \end{pmatrix} v^n$$

- In higher dimensions one first Fourier transforms in transversal variables. Eg. in 2D.

$$\begin{cases} \frac{d u_{i,j}}{dt} = Q_x u_{i,j} + Q_y u_{i,j} \\ u_{0,j} - 2u_{1,j} + u_{2,j} = 0 \\ u_0 = f_j \end{cases}$$



$$\text{Let } \hat{u}_i(\omega) = \sum u_{i,j} e^{ij\omega y} \quad (\text{periodic in } y)$$

$$\begin{cases} \frac{d \hat{u}_i}{dt} = (\hat{Q}_x + \hat{Q}_y) \hat{u}_i \\ \hat{u}_0 - 2\hat{u}_1 + \hat{u}_2 = 0 \end{cases}$$

\Rightarrow One ~~step~~-dimensional problem for each ω_y .

All conditions must now hold uniformly in ω_y .

(9)

For scalar problems

$$y_j^{n+1} = Q y_j^n$$

$$Q = \sum_{m=r}^p a_m E^m$$

$$y_j^0 = 0$$

$$v_{j\mu}^{n+1} = g_{j\mu}^{n+1}$$

$$v = -r+1, \dots, 0$$

always satisfies the Kreiss condition. (Goldberg & Tudmor)

Generalized stability.

Much theory developed via another stability concept: "strongly stable in generalized sense"

$$\int_0^{\infty} e^{-\eta t} \|u(t)\|^2 dt \leq K(\eta) \int_0^{\infty} e^{-\eta t} (\|F(t)\|^2 + |g(t)|^2) dt$$

when initial data $f=0$. Natural for Laplace / Z-transform setting.

- Enough to prove convergence (except if initial error very rough.)
- Schemes can be stable in this sense without satisfying the Kreiss cond. or being boundary stable.

a Strip problem

When we have two boundaries $0 < x < 1$ we have in general that if each boundary satisfies Kreiss condition and we have a energy estimate then strip problem is stable (in same sense as the Y_4 -problems are).

Variable coefficients

Stability for ^{all} Frozen coefficient problems

typically give stability also for variable coefficient problem. Hard to find precise statements of this though.